

# ***Turbulence, Erratic Property and Horseshoes in a Coupled Lattice System related with Belusov–Zhabotinsky Reaction***

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## **ABSTRACT**

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In this paper we continue to study the chaotic properties of the following lattice dynamical system  $b_j^{i+1} = \alpha_1 g(b_{j-1}^i) + \alpha_2 g(b_j^{i+1}) + \alpha_3 g(b_{j+1}^i)$ , where  $i$  is discrete time index,  $j$  is lattice side index with system size  $L$ ,  $g$  is a selfmap on  $[0, 1]$  and  $\alpha_1, \alpha_2, \alpha_3 \in [0, 1]$  with  $\alpha_1 + \alpha_2 + \alpha_3 = 1$  are coupling constants. In particular, it is shown that if  $g$  is turbulent (resp. erratic) then so is the above system, and that if there exists a  $g$ -connected family  $G$  with respect to disjointed compact subsets  $D_1, D_2, \dots, D_m$ , then there is a compact invariant set  $K' \subset D'$  such that  $F|_{K'}$  is semi-conjugate to  $m$ -shift for any coupling constants  $\alpha_1, \alpha_2, \alpha_3 \in [0, 1]$  with  $\alpha_1 + \alpha_2 + \alpha_3 = 1$ , where  $D' \subset I^L$  is nonempty and compact. Moreover, an example and two problems are given.

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## **1. INTRODUCTION AND PRELIMINARIES**

A topological dynamical system  $(G, g)$  is a compact metric space  $G$  together with a continuous map  $g: G \rightarrow G$  on  $G$ . Since Li and Yorke [1] first defined chaos in 1975, many dynamical properties in topological dynamical systems were highly discussed in the literatures (see [2–3]). For the importance of the lattice dynamical systems, we refer the reader to [4]. Many authors (see [5–13]) studied the following lattice dynamical system:

$$b_j^{i+1} = \alpha_1 g(b_{j-1}^i) + \alpha_2 g(b_j^{i+1}) + \alpha_3 g(b_{j+1}^i),$$

where  $i$  is discrete time index,  $j$  is lattice side index with system size  $L$ ,  $\alpha_1, \alpha_2, \alpha_3 \in [0, 1]$  with  $\alpha_1 + \alpha_2 + \alpha_3 = 1$ , are coupling constants and  $g$  is a continuous self map on  $I = [0, 1]$ .

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To analyze whether such a system is complicated or not by the observation of some topological dynamic property of this system is an open problem (see [5]). In [5], the authors characterized the dynamical complexity of a coupled lattice system stated by Kaneko in [14] which is related to the Belousov–Zhabotinsky reaction and deduced that this kind of system is Li–Yorke chaotic and Devaney chaotic for the case of  $\alpha_1 = 1$  and  $\alpha_2 = \alpha_3 = 0$ . Also, some problems on the dynamics of the system (1) with  $\alpha_1 < 1$  and  $\alpha_2 > 0$  (resp.  $\alpha_3 > 0$ ) are presented by them. Recently, in [13] the authors established that for  $\alpha_2 = \alpha_3 = \frac{1}{2}\alpha > 0$  and  $\alpha_1 = 1 - \alpha$  and the unimodal map  $g$  on  $I$ , this system (1) is Li–Yorke chaotic and has positive entropy.

The notion of distributional chaos was first introduced by Schweizer and Smítal in [15]. It is very important, because it is equivalent to the concept of positive topological entropy and some other kinds of chaos for compact intervals [15] or hyperbolic symbolic spaces [16]. But this equivalence is not true for higher dimensional spaces [17], the same happens if the dimension is zero [18]). In [19] the authors constructed a distributional chaotic minimal system. More recently, in [20] Wu and Zhu established that for  $\alpha_2 = \alpha_3 = 1/2\alpha > 0$  and  $\alpha_1 = 1 - \alpha$  and the tent map  $\Lambda$  defined by  $\Lambda(y) = 1 - |1 - 2y|$ , the system (1) is distributionally  $(a, b)$ -chaotic for any  $0 \leq a \leq b \leq 1$  and obtained that principal measure of such a system is not less than

$$\frac{2}{3} + \sum_{j=2}^{\infty} \frac{1}{j} \frac{2^{j-1}}{(2^j+1)(2^{j-1}+1)}$$

for the map  $\Lambda$ . Inspired by the above results, we will continue to consider some new chaotic properties of the system (1). In particular, we prove that for the system (1) and any  $\alpha_1, \alpha_2, \alpha_3 \in I$  with  $\alpha_1 + \alpha_2 + \alpha_3 = 1$ , if  $g$  turbulent or erratic, then so is the system (1). Moreover, we present an example shows that there is a turbulent and erratic coupled lattice system related with Belousov–Zhabotinsky reaction. For some problems, some properties and notions of dynamical systems, we refer the reader to [21].

Let  $(G, \rho)$  be a metric space and  $(G, g)$  be a dynamical system. A pair  $(a, b) \in G \times G$  is a Li–Yorke pair of the system  $(G, g)$  if

$$\limsup_{j \rightarrow \infty} \rho(g^j(a), g^j(b)) > 0$$

and

$$\liminf_{j \rightarrow \infty} \rho(g^j(a), g^j(b)) = 0$$

A subset  $E \subset G$  is a Li–Yorke set of  $g$  if any pair of distinct points in  $E$  is a Li–Yorke pair of the system  $(G, g)$ . A dynamical system  $(G, g)$  or a map  $g: G \rightarrow G$  is Li–Yorke chaotic if the space  $G$  contains an uncountable Li–Yorke set of  $g$ .

For any topological dynamical system  $(G, g)$ , any  $a, b \in G$  and any integer  $j > 0$ , the distributional function  $\Phi_{ab}^j: \mathbb{R}^+ \rightarrow [0, 1]$  is given by

$$\Phi_{ab}^j(t) = \frac{1}{j} \# \{i \in \mathbb{N}: \rho(g^i(a), g^i(b)) < t, 1 \leq i \leq j\},$$

where  $\mathbb{R}^+ = [0, +\infty)$  and  $\# A$  is the cardinality of the set  $A$ . Set

$$\Phi_{ab}(t, g) = \liminf_{j \rightarrow \infty} \Phi_{ab}^j(t)$$

and

$$\Phi_{ab}^*(t, g) = \limsup_{j \rightarrow \infty} \Phi_{ab}^j(t)$$

For any  $(a, c) \in I$  with  $a \leq b$ , a topological dynamical system  $(G, g)$  or a selfmap  $g: G \rightarrow G$  is distributionally  $(a, b)$ -chaotic if the space  $G$  contains an uncountable subset  $E$  such that for some  $\tau > 0$ ,  $\Phi_{xy}(t, g) = a$  and  $\Phi_{xy}^*(t, g) = b$  for any  $x, y \in E$  with  $x \neq y$  and any  $t \in (0, \tau)$ . Clearly, if the system  $(G, g)$  or the map  $g$  is distributionally  $(0, 1)$ -chaotic then it is distributional chaotic (see [20, 22]).

The principal measure  $\mu_p(g)$  of a topological dynamical system  $(G, g)$  or a selfmap  $g: G \rightarrow G$  is given by

$$\mu_p(g) = \sup_{a, b \in G} \frac{1}{l} \int_0^{+\infty} (\Phi_{ab}^*(t, g) - \Phi_{ab}(t, g)) dt$$

where  $l = \text{diam}(G)$  denotes the diameter of the space  $G$  (see [23]). From [23] one can know that

$$\mu_p(\Lambda) = \frac{2}{3} + \sum_{j=2}^{\infty} \frac{1}{j} \frac{2^{j-1}}{(2^{j+1})(2^{j-1}+1)}$$

where  $\Lambda$  is the tent map given by  $\Lambda(y) = 1 - |1 - 2y|$  for any  $y \in I$ .

Let  $g: G \rightarrow G$  be a continuous map on a topological space  $G$ . The map  $g$  is called a turbulent map (see [24]) if there are two nonempty closed subsets  $J, K \subset G$  with  $J \cap K = \emptyset$  such that  $J \cup K \subset g(J) \cap g(K)$ . The map  $g$  is said to be erratic (see [24]) if there is a nonempty closed set  $B \subset G$  such that  $B \cap g(B) = \emptyset$  and  $B \cup g(B) \subset g^2(B)$ .

The state space of the system (1) is the set

$$G = \{\mathbf{b}: \mathbf{b} = \{\mathbf{b}_j\}, \mathbf{b}_j \in \mathbb{R}^p, \mathbf{j} \in \mathbb{Z}^q, \|\mathbf{b}_j\| < \infty\},$$

where  $p \geq 1$  is the dimension of the range space of the map of state  $\mathbf{b}_j$ ,  $|\mathbf{b}_j|$  is the length of the vector  $\mathbf{b}_j$ ,  $q \geq 1$  is the dimension of the lattice and the  $l^2$  norm

$$\|\mathbf{b}\|_2 = \left(\sum_{\mathbf{j} \in \mathbb{Z}^q} |\mathbf{b}_i|^2\right)^{\frac{1}{2}} \tag{1}$$

is usually taken (see [5]). In general, one of the following periodic boundary conditions of the system (1) is needed:

1.  $b_j^i = b_{j+L}^i$ ,
2.  $b_j^i = b_j^{j+L}$ ,
3.  $b_j^i = b_{j+L}^{i+L}$ ,

where standardly, the first case of the boundary conditions is used.

## 2. MAIN RESULTS

The system (1) with  $\alpha_1 = 1 - \alpha$  and  $\alpha_2 = \alpha_3 = \alpha/2$  and  $\alpha \in I$  was explored by lots of authors, mostly experimentally or semi-analytically than analytically. The first paper having analytic results is [25], where the authors got that the system (1) with  $\alpha_1 = 1 - \alpha$ ,  $\alpha_2 = \alpha_3 = \alpha/2$  and  $\alpha \in I$  is Li-Yorke chaotic. In [5] the authors presented an alternative and easier proof of the above result and claimed that the system (1) with  $\alpha_1 = 1 - \alpha < 1$ ,  $\alpha_2 = \alpha_3 = \alpha/2 > 0$  and  $\alpha \in I$  is more complicated.

Let  $\tilde{\rho}$  be the product metric on the product space  $I^L$ , i.e.,

$$\tilde{\rho}((x_1, x_2, \dots, x_L), (y_1, y_2, \dots, y_L)) = (\sum_{i=1}^L (x_i - y_i)^2)^{\frac{1}{2}}$$

for any  $(x_1, x_2, \dots, x_L), (y_1, y_2, \dots, y_L) \in I^L$ . Define the map  $F: (I^L, d) \rightarrow (I^L, d)$  by  $F(x_1, x_2, \dots, x_L) = (y_1, y_2, \dots, y_L)$  where  $y_i = \alpha_1 g(x_i) + \alpha_2 g(x_{i-1}) + \alpha_3 g(x_{i+1})$ ,  $\alpha_1, \alpha_2, \alpha_3 \in I$  and  $\alpha_1 + \alpha_2 + \alpha_3 = 1$ . It is obvious that the system (1) is equivalent to the system  $(I^L, F)$ . In [20] the authors got that if  $g = \Lambda$ , then the system (1) with  $\alpha_1 = 1 - \alpha$ ,  $\alpha_2 = \alpha_3 = \alpha/2$  and  $\alpha \in I$  is distributionally  $(a, b)$ -chaotic for any  $(a, b) \in I \times I$  with  $a \leq b$  and any  $\alpha \in (0, 1)$ . Inspired by the above results we have the following theorems.

**Theorem 2.1.** For any coupling constants  $\alpha_1, \alpha_2, \alpha_3 \in [0, 1]$  with  $\alpha_1 + \alpha_2 + \alpha_3 = 1$  and any selfmap  $g$  on  $[0, 1]$ , if  $g$  is turbulent, then so is the system (1).

**Proof.** Fix  $\alpha_1, \alpha_2, \alpha_3 \in [0, 1]$  with  $\alpha_1 + \alpha_2 + \alpha_3 = 1$ . By hypothesis and the definition, there are two nonempty closed subsets  $J, K \subset G$  with  $J \cap K = \emptyset$  such that  $J \cup K \subset g(J) \cap g(K)$ . Let

$$J' = \{(x, x, \dots, x) \in I^L : x \in J\}$$

and

$$K' = \{(y, y, \dots, y) \in I^L : y \in K\}.$$

Then we have that  $J', K' \subset I^L$  are two nonempty closed subsets with  $J' \cap K' = \emptyset$  and

$$J' \cup K' \subset \underbrace{g \times g \times \dots \times g}_L(J') \cap \underbrace{g \times g \times \dots \times g}_L(K').$$

By the definition, the system (1) is turbulent.

**Problem 2.1.** Let  $\alpha_1, \alpha_2, \alpha_3 \in [0, 1]$  with  $\alpha_1 + \alpha_2 + \alpha_3 = 1$  and the system (1) be turbulent. Is  $g$  turbulent?

**Theorem 2.2.** For any coupling constants  $\alpha_1, \alpha_2, \alpha_3 \in [0, 1]$  with  $\alpha_1 + \alpha_2 + \alpha_3 = 1$  and any selfmap  $g$  on  $[0, 1]$ , if  $g$  is erratic, then so is the system (1).

**Proof.** Fix  $\alpha_1, \alpha_2, \alpha_3 \in [0, 1]$  with  $\alpha_1 + \alpha_2 + \alpha_3 = 1$ . By hypothesis and the definition, there is a nonempty closed subset  $J \subset G$  with  $J \cap g(J) = \emptyset$  such that  $J \cup g(J) \subset g_2(J)$ . Let  $J' = \{(x, x, \dots, x) \in I^L : x \in J\}$ . Then we have that  $J' \subset I^L$  is a nonempty closed subset with  $J' \cap \underbrace{g \times g \times \dots \times g}_L(J') = \emptyset$  and

$$J' \cup \underbrace{g \times g \times \dots \times g}_L(J') \subset \underbrace{(g \times g \times \dots \times g)}_L^2(J').$$

By definition, the system (1) is erratic.

**Problem 2.2.** Let  $\alpha_1, \alpha_2, \alpha_3 \in [0, 1]$  with  $\alpha_1 + \alpha_2 + \alpha_3 = 1$  and the system (1) be erratic. Is  $g$  erratic?

**Example 2.1.** Let  $g$  be the tent map defined by  $g(y) = 1 - |1 - 2y|$ ,  $y \in [0, 1]$ . Then  $g$  is erratic and turbulent. Consequently, for the tent map and any coupling constants  $\alpha_1, \alpha_2, \alpha_3 \in [0, 1]$  with  $\alpha_1 + \alpha_2 + \alpha_3 = 1$ , the system (1) is erratic and turbulent.

**Proof.** From Example 1 in [24] we know that the tent map  $g$  is erratic. By the definition and Example 1 in [24] one can see that the tent map  $g$  is turbulent. By Theorems 2.1 and 2.2, the system (1) is erratic and turbulent.

Now we recall some aspects of symbolic dynamics (see [26–28]). Let  $S_m = \{0, 1, \dots, m-1\}$ ,  $Z_+ = \{0, 1, \dots\}$ ,  $\Sigma_m^+ = \Sigma_+^{S_m}$  and  $\Sigma_m = \Sigma^{S_m}$ . For any sequences  $s = \{\dots, s_{-n}, \dots, s_{-1}, s_0, s_1, \dots, s_n, \dots\}$ ,  $\bar{s} = \{\dots, \bar{s}_{-n}, \dots, \bar{s}_{-1}, \bar{s}_0, \bar{s}_1, \dots, \bar{s}_n, \dots\} \in \Sigma_m$  (or  $s = \{s_0, s_1, \dots, s_n, \dots\}$ ,  $\bar{s} = \{\bar{s}_0, \bar{s}_1, \dots, \bar{s}_n, \dots\} \in \Sigma_m^+$ ), the distance between  $s$  and  $\bar{s}$  is defined by  $d(s, \bar{s}) = \sum_{i=-\infty}^{\infty} \frac{1}{2^{|i|}} \frac{|s_i - \bar{s}_i|}{1 + |s_i - \bar{s}_i|}$  in case of bi-infinite sequences (or  $d(s, \bar{s}) = \sum_{i=0}^{\infty} \frac{1}{2^{|i|}} \frac{|s_i - \bar{s}_i|}{1 + |s_i - \bar{s}_i|}$  in case of one-sided sequences). With the distance defined as above,  $\Sigma_m$  (or  $\Sigma_m^+$ ) is a perfect, totally disconnected and compact metric space (see [26–27]). A  $m$ -shift map  $\sigma: \Sigma_m \rightarrow \Sigma_m$  (or  $\sigma: \Sigma_m^+ \rightarrow \Sigma_m^+$ ) is defined by  $\sigma(s)_i = s_{i+1}$  for any  $s \in \Sigma_m$  (or  $s \in \Sigma_m^+$ ).

**Definition 2.1.** (see [26–27]) Let  $X$  be a metric space and  $g: X \rightarrow X$  be a continuous map. Let  $K \subset X$  be a compact invariant set of  $g$ . If there is a continuous surjective map  $h: K \rightarrow \Sigma_m$  (or  $h: K \rightarrow \Sigma_m^+$ ) such that  $h \circ g = \sigma \circ h$ , then  $g|_K$  is said to be semi-conjugate to  $\sigma$ . Let  $X$  be a metric space,  $D \subset X$  be compact and  $D_1, D_2, \dots, D_m$  be mutually disjoint compact subsets of  $D$ . Suppose that  $g: D_i \rightarrow X$  is a continuous map for every  $i \in \{1, 2, \dots, m\}$ .

**Definition 2.2.** (see [26–27]) Let  $D_1, D_2, \dots, D_m$  be mutually disjoint compact subsets of  $D$ . Assume that  $\gamma$  is a compact connected subset of  $D$  such that for each  $i \in \{1, 2, \dots, m\}$ ,  $\gamma_i = \gamma \cap D_i$  is nonempty and compact. Then we say that  $\gamma$  is a connection with respect to  $D_1, D_2, \dots, D_m$ . Let  $G$  be a family of connections  $\gamma_s$  with respect to  $D_1, D_2, \dots, D_m$  such that  $\gamma \in G$  implies  $g(\gamma_i) \in G$  for each  $i \in \{1, 2, \dots, m\}$ . Then we say that  $G$  is a  $g$ -connected family with respect to  $D_1, D_2, \dots, D_m$ .

**Horseshoe Lemma** (see [26–27]) Suppose that there exists a  $g$ -connected family  $G$  with respect to disjoint compact subsets  $D_1, D_2, \dots, D_m$ . Then there is a compact invariant set  $K \subset D$  such that  $g|_K$  is semi-conjugate to  $m$ -shift. Now we establish the following result.

**Theorem 2.3.** For any coupling constants  $\alpha_1, \alpha_2, \alpha_3 \in [0, 1]$  with  $\alpha_1 + \alpha_2 + \alpha_3 = 1$ , if there exists a  $g$ -connected family  $G$  with respect to disjoint compact subsets  $D_1, D_2, \dots, D_m$  then, for the system (1) there is a compact invariant set  $K' \subset D'$  such that  $F|_{K'}$  is semiconjugate to  $m$ -shift, where  $D' \subset I^L$  is nonempty and compact.

**Proof.** Write  $\Delta_{[0,1]^L} = \{(x_1, x_2, \dots, x_L) : x_1 = x_2 = \dots = x_L \in [0, 1]\}$ . Since  $F$  is a continuous selfmap of  $I^L$ ,  $(\Delta_{[0,1]^L}, F|_{\Delta_{[0,1]^L}})$  is a subsystem of the system  $(I^L, F)$ . Define a map  $h: \Delta_{[0,1]^L} \rightarrow [0, 1]$  by  $h(\vec{a}) = a$  for any  $\vec{a} = (a, a, \dots, a) \in \Delta_{[0,1]^L}$ . It is easily verified that  $h$  is a homeomorphism. Clearly, we have  $h \circ F|_{\Delta_{[0,1]^L}}(\vec{x}) = h(\overrightarrow{g(x)})$  and  $h(\overrightarrow{g(x)}) = g(x) = g \circ h(\vec{x})$ . So,  $h \circ F|_{\Delta_{[0,1]^L}} = g \circ h$ . This shows that  $(\Delta_{[0,1]^L}, F|_{\Delta_{[0,1]^L}})$  is topologically conjugate to the system  $([0, 1], g)$ . Clearly, horseshoe chaos is invariant under topological conjugation. Consequently, Theorem 3.1 holds. Thus, the proof is ended.

### 3. REMARK

Similar to the above discussion, one can prove that the results in this paper hold for the following coupled map lattice:

$$b_{m+1,n} = \alpha_{-s}g(b_{m,n-s}) + \alpha_{-s+1}g(b_{m,n-s+1}) + \dots + \alpha_0g(b_{m,n}) + \dots + \alpha_tg(b_{m,n+t}) \quad (2)$$

where  $s, t \in \mathbb{Z}$ ,  $s, t \geq 0$ ,  $\alpha_{-s}, \dots, \alpha_0, \dots, \alpha_t > 0$ ,  $\alpha_{-s} + \dots + \alpha_0 + \dots + \alpha_t = 1$ .

In fact, by  $\alpha_{-s} + \dots + \alpha_0 + \dots + \alpha_t = 1$ ,  $F(\vec{b}) = \overrightarrow{g(b)}$ , for any  $\vec{b} \in I^L$ . This implies that the the system (2) has the same dynamical properties as the system (1).

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